

# The Topos Approach

in the Qubit case

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Heunen/Landsman/  
Spitters

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# The topos approach to quantum formalism

Describe quantum system by  
von Neumann algebra

A  $\left\{ \begin{array}{l} \text{von Neumann algebra} \\ \text{C}^* \text{-algebra} \end{array} \right.$

Isham/Butterfield/  
Döring

Heunen/Landsman/  
Spitters

Observable = self-adjoint element of A

in  $A_{sa}$

## Classical case - A commutative

A has a (Gelfand-Naimark) spectrum  $\Sigma$ ,

$$A \cong C(\Sigma, \mathbb{C}) \quad A_{sa} \cong C(\Sigma, \mathbb{R})$$

Element  $\phi \in \Sigma$  is classically pure state -

specifies value of every observable

Also -  $\text{Dist}(\Sigma) =$  space of distributions on  $\Sigma$  regular probability measures

$\phi \in \text{Dist}(\Sigma)$  is mixed state - specifies probabilistic distribution of values of observables

## Quantum case - A non-commutative

e.g. A = bdd linear operators on Hilbert space  $\mathcal{H}$ .

$\phi \in \mathcal{H}$  gives probabilistic distribution of values of observables Born rule

BUT no spectrum of classically pure states

Kochen-Specker:

Cannot assign values to all observables consistent with their functional relationships

$$\begin{array}{l} a = f(b) \\ f: \mathbb{R} \rightarrow \mathbb{R} \end{array}$$

## Topos approach

- Define a topos  $\mathcal{J}(A)$
- Define spectrum in **internal mathematics** of  $\mathcal{J}(A)$
- Do internal mathematics, but extract external information.

What does this mean?

Let  $\mathcal{L}(A) =$  posets of commutative subalgebras of  $A$   
— "classical viewpoints" of system

Each  $C \in \mathcal{L}(A)$  has spectrum  $\Sigma(C)$

Idea: classical treatment  
indexed by classical viewpoint  
→ quantum treatment

## Topos internal as bundles

Suppose topos  $\mathcal{J} = \text{Sh}(X)$

Object of  $\mathcal{J} =$  local homeomorphism over  $X$

Internal space in  $\mathcal{J}$  **use point-free topology**

= bundle over  $X$

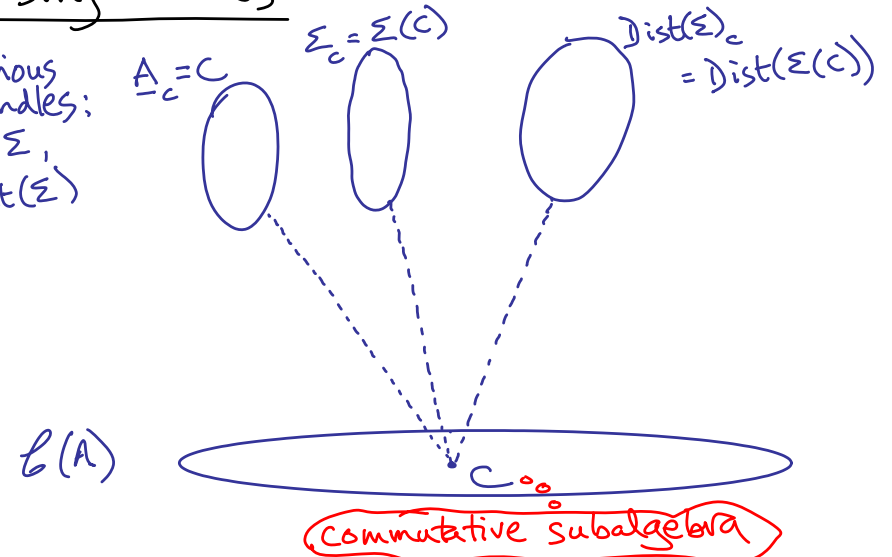
Topos-valid, point-free reasoning internally

= bundles externally **geometricity**

often — internal reasoning  
= fibrewise, indexed by base point

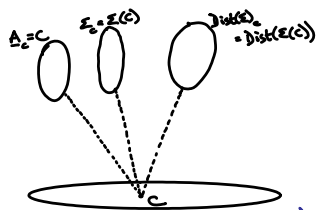
Using bundles Matches Landsman et al.

Various bundles:  
 $A, \Sigma, \text{Dist}(\Sigma)$



Kochen-Specker

$\Rightarrow \Sigma$  has no cross-sections  
(continuous choice of point  
from every fibre = external point of bundle)



However —  $\text{Dist}(\Sigma)$  does have cross-sections

Each quantum state produces one.

Internally: quantum states are distributions  
over classical states

Externally: this is impossible

Does topos support "neo realist" reasoning?  
oo Isham

Isham/Döring:  $\mathcal{Y}(A) = \text{Set}^{\mathcal{C}(A)^{\text{op}}} = \text{Sh}(\text{Filt}(\mathcal{C}(A)))$   
filter completion

Heunen/Landsman/Spitters:  $\mathcal{Y}(A) = \text{Set}^{\mathcal{C}(A)} = \text{Sh}(\text{Idl}(\mathcal{C}(A)))$   
ideal completion

Both: — include classical viewpoints as  
base points (principal filters/ideals)  
— respect order on  $\mathcal{C}(A)$  but not topology  
Idea: use topology on  $\mathcal{C}(A)$

Case  $A = M_2(\mathbb{C})$

Let  $\mathcal{C}_2(A) = \{C \in \mathcal{C}(A) \mid \dim C = 2\}$  greatest dimension possible

$C \in \mathcal{C}_2(A) \Rightarrow C \cong \mathbb{C}^2$ , spectrum = 2

$C$  determined by two projectors  $E, I-E$

cor. to  $(1,0), (0,1)$

Spectrum =  $\{E, I-E\}$

self-adjoint idempotents

Projector has — eigenvalues 0 or 1

$\therefore$  trace 0, 1 or 2

trace 0, 2 are  $\begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$  &  $\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$

We seek projectors  $E$  with trace 1.

$E, I-E$  both describe same  $C$ .

Projectors via unitaries

$E$  idempotent  $\Leftrightarrow U = 2E - I$  has  $U^2 = I$

$E$  a projector  $\Leftrightarrow U$  s.a. unitary

$\text{tr} E = 1 \Leftrightarrow \text{tr} U = 0$

s.a. unitaries, trace 0:

$U = \begin{pmatrix} c & a+ib \\ a-ib & -c \end{pmatrix}$ ,  $a^2 + b^2 + c^2 = 1$

$E, I-E$  cor. to  $U, -U$

$\mathcal{C}_2(A) \cong$  sphere  $S^2$  with antipodes identified  
= real projective plane  $\mathbb{R}P^2$

Take  $\mathcal{Y}(A) = \text{Sh}(\mathcal{C}_2(A)) = \text{Sh}(\mathbb{R}P^2)$

Spectral bundle is obvious cover  $S^2$

- A local homeomorphism  
(object of  $\mathcal{Y}(A)$ )

- Each fibre has two points

- continuous No global sections

- cf. Kochen-Specker (doesn't apply to  $M_2(A)$ , but doesn't require continuity)

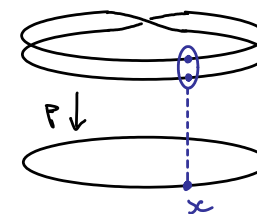
$\downarrow$   
 $\mathbb{R}P^2$

No global sections

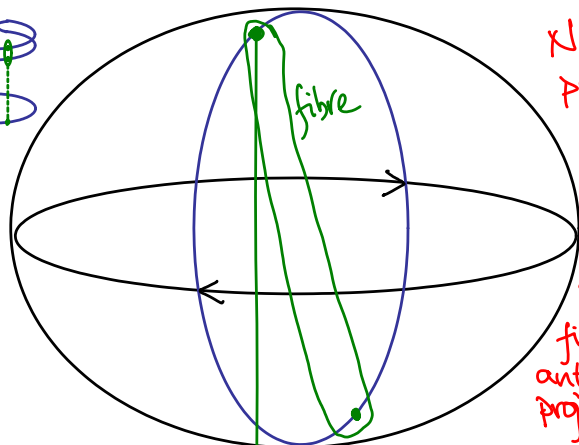
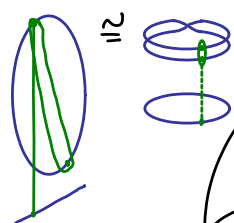
$S^2 \cong S^1$

$\downarrow$

$\mathbb{R}P^2 \cong S^1$



Global section would give continuous choice of edge of Möbius strip.



N. hemisphere projects down to disc

S. hemisphere: first take antipode, then project down

Projective plane  
points at  $\infty$

Disc, identity ends of diameters

Distributions

point-free  
ie. locale

Internally: space  $X \mapsto$  space  $\text{Dist}(X)$   
of distributions (regular probability measures)

Externally: bundle  $\mapsto$  distribution bundle  
*Heunen/Landsman/Spitzer*

Logical argument  $\Rightarrow$  construction works fibrewise  
geometricity of Distribution big result!

## Global distributions

General feature of topos approach

Internally: have spectrum  $\Sigma$

element = classically pure state

BUT no global elements - can't specify results of all observables simultaneously

Also have  $\text{Dist}(\Sigma)$

element = classically mixed state

DES have global elements - e.g. each quantum pure state provides one.

(External) quantum pure = classically mixed - but only internally.

e.g. for  $A = M_2(\mathbb{C})$

On fibres:

$$\text{Dist}(2) = \{(p_0, p_1) \in [0,1] \mid p_0 + p_1 = 1\} \cong [0,1]$$

probabilities of two points of spectrum

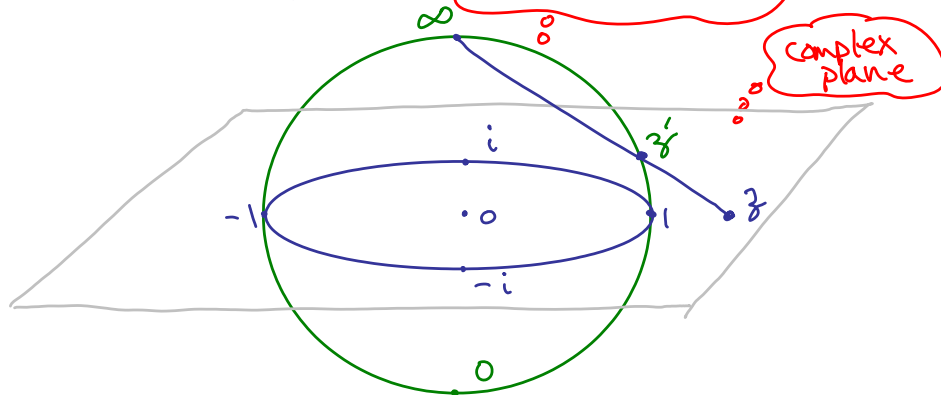
Show how quantum pure state  $\in \mathbb{C}^2$

gives global section of  $\text{Dist} \left( \begin{array}{c} S^2 \\ \downarrow \\ \mathbb{R}P^2 \end{array} \right)$

## States $\phi \in \mathbb{C}^2$

Scale invariant,  $\therefore$  in complex projective line  $\mathbb{C}P^1$

$\cong S^2$  Riemann sphere



If  $z' = (a, b, c)$  on sphere

$$a^2 + b^2 + c^2 = 1$$

$$c \neq 1$$

then  $z = \frac{a}{1-c} + i \frac{b}{1-c}$  on plane

$$(a, b, c) = c(0, 0, 1) + (1-c)\left(\frac{a}{1-c}, \frac{b}{1-c}, 0\right)$$

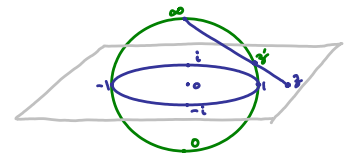
- corresponds to 1-dim subspace

$$\mathbb{C} \begin{pmatrix} a+ib \\ 1-c \end{pmatrix} \text{ in } \mathbb{C}^2$$

= eigenspace for  $E(a, b, c)$

with eigenvalue 1

$$\begin{aligned} & \frac{1}{2} \begin{pmatrix} 1+c & a+ib \\ a-ib & 1-c \end{pmatrix} \begin{pmatrix} a+ib \\ 1-c \end{pmatrix} \\ &= \frac{1}{2} (a+ib)(1+c+1-c) \\ &= \frac{1}{2} (a^2+b^2+1-2c+c^2) \\ &= \begin{pmatrix} a+ib \\ 1-c \end{pmatrix} \end{aligned}$$



Consider: projector  $\begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$  - North Pole  
 $E(0,0,1)$

state  $\phi = \begin{pmatrix} a+ib \\ 1-c \end{pmatrix}$

Born rule: probability of going to  $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$   
= expected value of  $E$

$$\begin{aligned} &= \frac{(a-ib, 1-c) \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} a+ib \\ 1-c \end{pmatrix}}{(a-ib, 1-c) \begin{pmatrix} a+ib \\ 1-c \end{pmatrix}} = \frac{(a-ib, 1-c) \begin{pmatrix} a+ib \\ 0 \end{pmatrix}}{a^2+b^2+1-2c+c^2} \\ &= \frac{a^2+b^2}{2(1-c)} = \frac{1+c}{2} \end{aligned}$$

In general

$E, \phi$  two points on  $S^2$   
projector state

Let  $P$  = equatorial plane taking  $E$  as North pole

Expected value  $p_E$  of  $E$  at  $\phi = \frac{1+c}{2}$  where  
 $c$  = distance of  $\phi$  from  $P$ , taking  $E$  to be on positive side.

Continuous, & coordinate independent

Fix state  $\phi$

For each  $C \in \mathbb{R}P^2$  -

- let  $E, E'$  be antipodal points in its fibre in  $S^2$
- $p_E + p_{E'} = 1$  (both calculated for  $\phi$ )
- Get point in fibre over  $C$  of distribution bundle
- Gives continuous global section of distribution bundle.

Product distributions

$\text{Dist}(X) \times \text{Dist}(Y) \rightarrow \text{Dist}(X \times Y)$

e.g.  $X, Y$  finite discrete

$$(p_x)_{x \in X} \quad (q_y)_{y \in Y} \mapsto (p_x q_y)_{\substack{x \in X \\ y \in Y}}$$

$$\sum_x p_x = 1, \quad \sum_y q_y = 1 \Rightarrow \sum_{x,y} p_x q_y = 1$$

Conjecture:  
still works internally

Not all distributions on product are product distributions

Distributions on —

- X -  $X \times 1$  column vector  $p$
  - $\gamma$  -  $1 \times \gamma$  row vector  $q$
  - $X \times \gamma$  -  $X \times \gamma$  matrix  $M$
- } with elements in  $[0,1]$  and summing to 1

Product distribution for  $p, q$  has matrix  $pq$

- rank = 1

$\therefore \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}, \begin{pmatrix} 0 & \frac{1}{2} \\ \frac{1}{2} & 0 \end{pmatrix}$  not product distributions - exhibit "entanglement" between  $X, \gamma$  components.

Product spaces

For  $\mathcal{A} \otimes \mathcal{A}'$  with  $A \otimes A'$ :

- don't know all commutative subalgebras
- if  $C, C' \in \mathcal{C}(A), \mathcal{C}(A')$  then  $C \otimes C' \in \mathcal{C}(A \otimes A')$
- what can we get from  $\sum x_i e_i$  ?  
 $\downarrow$   
 $\mathcal{C}(A) \times \mathcal{C}(A')$

$\mathcal{A} = \mathcal{A}' = \mathbb{C}^2$

Product states  $\phi \phi'$  ( $= \phi \otimes \phi'$ ),  $\phi, \phi'$  unit vectors

If  $C$  given by projectors  $E, I-E$

...  $C'$  ...  $E', I-E'$

$\phi \phi'$  gives distribution matrix

$$\begin{pmatrix} \langle \phi \phi', (E \otimes E') \phi \phi' \rangle & \langle \phi \phi', (E \otimes (I-E')) \phi \phi' \rangle \\ \langle \phi \phi', ((I-E) \otimes E') \phi \phi' \rangle & \langle \phi \phi', ((I-E) \otimes (I-E')) \phi \phi' \rangle \end{pmatrix}$$

$= \begin{pmatrix} \langle \phi, E \phi \rangle \\ \langle \phi, (I-E) \phi \rangle \end{pmatrix} \begin{pmatrix} \langle \phi', E' \phi' \rangle, \langle \phi', (I-E') \phi' \rangle \end{pmatrix}$  - product distribution

Entangled state? e.g.  $\psi = (\uparrow\uparrow + \downarrow\downarrow)/\sqrt{2}$   $\uparrow = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$   
 $\downarrow = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$

If  $E = \frac{1}{2} \begin{pmatrix} 1+c & a+ib \\ a-ib & 1-c \end{pmatrix}$

need  $\langle \uparrow, E \uparrow \rangle = \frac{1+c}{2}$

$\langle \downarrow, E \downarrow \rangle = \frac{1-c}{2}$

$\langle \uparrow, E \downarrow \rangle = \frac{a-ib}{2}$

$\langle \downarrow, E \uparrow \rangle = \frac{a+ib}{2}$

calculated before

$$\langle \psi, (E \otimes E') \psi \rangle = \frac{1 + cc' + aa' - bb'}{4}$$

coordinate dependent?

Four projector pairs - take antipodes

$$\text{Distribution } \frac{1}{4} \begin{pmatrix} 1+cc'+aa'-bb' & 1-cc'-aa'+bb' \\ 1-cc'-aa'+bb' & 1+cc'+aa'-bb' \end{pmatrix}$$

e.g. for  $c, c' = \pm 1$

$$\text{Distribution } \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}$$

not a product - rank = 2

## Conclusions

- Illustrated external bundle view of topos internal reasoning
- Example using topology of  $G(A)$
- "Kochen-Specker" even in dimension 2, if require continuity
- Global distributions
- Some manifestation of entanglement.